Structural Investigation of Glasses Belonging to the Na₂O–Nb₂O₅–P₂O₅ System*

A. EL JAZOULI

Laboratoire de Chimie du Solide du CNRS, Université de Bordeaux I, 351 Cours de la Libération, 33405 Talence Cedex, France

J. C. VIALA

Laboratoire de Physico-Chimie Minérale, Université de Lyon I, 43 Boulevard du 11 Nov. 1918, 69621 Villeurbanne, France

AND C. PARENT, G. LE FLEM, AND P. HAGENMULLER

Laboratoire de Chimie du Solide du CNRS, Université de Bordeaux I, 351 Cours de la Libération, 33405 Talence Cedex, France

Received May 18, 1987

The structure of glasses [with composition 0.855 ($xNa_2O-yNb_2O_5-zP_2O_5$)- $0.145 Ln_2O_3$; Ln = La, Nd, Eu, and x + y + z = 100] has been investigated by Raman scattering and optical spectroscopy using Nd³⁺ and Eu³⁺ as local structural probes. For the niobium-poor glasses (y < 12.5) octahedral NbO₆ groups are inserted within modified metaphosphate chains. In contrast the network former of the niobium-rich glasses (y > 12.5) contains chains of NbO₆ octahedra sharing common corners. An intermediate situation seems to occur for $y \approx 12.5$ which preserves the Nasicon-type structure with all oxygen atoms common to NbO₆ and PO₄ groups. @ 1988 Academic Press, Inc.

I. Introduction

Due to their particular physical chemical properties, niobium oxide-based glasses appear quite attractive in various fields of material science. They can be used, for instance, for the elaboration of the $MNbO_3$ (M = Li, Na) glass ceramics (1-3) or as neodymium laser materials with a low optical pumping threshold and a very narrow emission bandwidth (4). Yet, such glasses have so far been poorly investigated. In a recent study the vitreous domains of the $M_2O-Nb_2O_5-P_2O_5$ (M = Li, Na) systems have been well-defined (5). They are quite similar whatever the alkali element considered, either Li or Na (Fig. 1). The lithium- and niobium-rich glasses exhibit good ionic conductivity which can reach $10^{-2} \Omega^{-1} \text{ cm}^{-1}$ at 300°C for the molar composition $50Li_2O-24Nb_2O_5-26P_2O_5$ with an activation energy $\Delta E = 0.45 \text{ eV}$.

To provide better insight into the structure of those glasses a detailed investigation was carried out by Raman scattering and optical spectroscopy using Eu³⁺ and Nd³⁺ as local structural probes. The present pa-

^{*} Dedicated to Professor A. Rabenau on the occasion of his 65th anniversary.



FIG. 1. Vitreous domain (hatched) in the Na₂O-Nb₂ O₅-P₂O₅ system. Glass compositions investigated for Na/P = 4/3(AB); C, y = 6; D, y = 12.5 (Na₄Nb(PO₄)₃); E, y = 18; F, y = 24.

per reports results obtained for rare-earth oxide-doped sodium glasses belonging to the Na₂O-Nb₂O₅-P₂O₅ system with Na/P = 1.33, i.e., on the AB line of Fig. 1. This line includes the compound Na₄Nb(PO₄)₃ for which a reversible glass-crystal transition was pointed out (6).

II. Experimental

As mentioned, the glasses investigated had molar compositions 0.855 (xNa_2O-y $Nb_2O_5-zP_2O_5$)-0.145 Ln_2O_3 with Ln = La, Nd, Eu, x + y + z = 100, and $x/z = \frac{4}{3}$. The y values were 6, 12.5, 18, and 24.

Optically inactive lanthanum was introduced for the Raman scattering and diffuse reflectance investigations in order to allow useful comparisons in terms of structural information for the lightly neodymium- or europium-doped glasses.

The glasses were prepared according to a procedure previously reported (5). Their thermal behavior was investigated using differential thermal analysis and classical X-ray diffraction (XRD) techniques. As an example, the vitreous transition and recrystallization temperatures of the rare-earth free glass corresponding to y = 12.5 were found to be 430 and 570°C, respectively. Recrystallization of the glasses led to Na₄ Nb(PO₄)₃, the XRD spectrum of which can be indexed assuming a monoclinic cell with parameters consistent with a Nasicon-type structure (6). For y = 6 the same phase is also found among the recrystallization products, whereas the XRD patterns corresponding to the recrystallized niobiumrich glasses (y = 18 and y = 24) are complex but don't include the characteristic lines of Na₄Nb(PO₄)₃.

The Raman spectra were recorded at room temperature using a Jobin-Yvon spectrophotometer coupled with an Ar laser Spectra Physics 2560. Samples in the form of pressed pellets ($\emptyset = 13$ mm, thickness = 1.5 mm) were prepared from ground glasses.

Reflectance diffuse spectra and absorption spectra of the neodymium-doped samples were determined at room temperature using a Cary 17 spectrophotometer.

The fluorescence of Nd³⁺ and Eu³⁺ was recorded using a Jobin-Yvon monochromator and a Varian VPM 159A photomultiplier. Steady-state emission spectra were obtained with a Xe lamp excitation. The lifetime of the Nd³⁺ ${}^{4}F_{3/2}$ level was deduced from the ${}^{4}F_{3/2} \rightarrow {}^{4}I_{11/2}$ fluorescence decay after a pulsed nonselective excitation in the ${}^{4}G_{5/2}$ and ${}^{2}G_{7/2}$ levels. The source was a rhodamine 590 solution excited by a nitrogen laser; the photomultiplier was connected to a PAR M 162 boxcar averager.

III. Raman Scattering Investigation

The Raman spectra of the investigated glasses are shown in Fig. 2. Band assignment can be made on the basis of previous data characterizing either sodium phosphate glasses (7, 8) or the Nb-O bonds in crystals (9-12).

For the composition close to the sodium phosphate (y = 6), the bands at 725 and 1145 cm⁻¹ can be assigned, respectively, to $\gamma_{s_{POP}}$ and $\gamma_{s_{PO_2}}$ vibrations of metaphosphate chains. Introduction of niobium gives rise to an intense band at 900 cm⁻¹ and to



FIG. 2a. Raman spectra of the 0.855 ($xNa_2O-yNb_2O_5-zP_2O_5$)-0.145 La₂O₃ glasses (x/z = 1.33; y = 6, 12.5, 18, and 24).

FIG. 2b. Comparison of the band intensities of the investigated glass spectra.

several weaker bands at 630, 450, and 260 cm^{-1} .

As the amount of niobium present in the glasses increases, the intensities of these new bands are enhanced, in particular that at 900 cm⁻¹. On the contrary the bands at 1145 and 725 cm⁻¹ weaken gradually, indicating a modification in the metaphosphate chains.

The new bands characterize Nb–O bonds in octahedral sites: the first two frequencies represent the Nb–O stretching mode, thus demonstrating the existence of a distribution of Nb–O lengths in the NbO₆ groups. The vibration at 900 cm⁻¹ corresponds to a short Nb–O bond ($d_{Nb-O} \approx 1.77$ Å), that at 630 cm⁻¹ is equivalent to the mode characterizing equatorial Nb–O bonds found in α -NbPO₅ ($d_{\rm Nb-O} \approx 1.97$ Å). In agreement with Ref. (11) the band at 450 cm⁻¹ probably involves O–Nb–O and O–P–O coupled modes. The remaining band located at 260 cm⁻¹ can be associated with various O– Nb–O bending modes. Such a band was also found in some ferroelectric materials with tungsten–bronze structures (12) in which the NbO₆ groups share common corners.

From those results it appears that the introduction of increasing amounts of Nb_2O_5 in the phosphate glasses initially breaks the metaphosphate chains and progressively develops a new network former structure constituted mainly by chains of NbO_6 octahedra sharing common corners.

An intermediate structure seems to occur for the glass corresponding to a formulation Na₄Nb(PO₄)₃ (y = 12.5), the crystallization of which leads to a Nasicon-type arrangement. Each NbO₆ unit is surrounded in this material by six PO₄ groups.

IV. Investigation of the Diffuse Reflectance Spectra

The diffuse reflectance spectra have been recorded and compared to those of crystalline materials. Figure 3 gives the spectra of the glasses as well as those of crystallized



FIG. 3. Comparison of the diffuse reflectance spectra of niobium crystallized oxides with those of the investigated glasses.

Absorption Threshold of $O^{2-} \rightarrow NO^{3+}$ Electronic Transfer in Niobium Crystallized Compounds and in the Investigated Glasses							
	0.855	Glass 0.855 (xNa ₂ O-yNb ₂ O ₅ -zP ₂ O ₅)-0.145 Ln ₂ O ₃					
	y = 6	y = 12.5	y = 18	<i>y</i> = 24	NbPO ₅	NaNbO ₃	
P/Nb	6.715	3	1.953	1.358	1	0	
Eu La	3.90 3.85	3.76 3.78	3.65 3.57	3.54 3.42	3.49	3.29	
	P/Nb Eu La	PTION THRESHOLD OF C COMPOL 0.855 (y) = 6 $y = 6$ $P/Nb 6.715$ $Eu 3.90$ $La 3.85$	PTION THRESHOLD OF $O^2 \rightarrow Nb^{5+}$ ELL COMPOUNDS AND IN TH Gl: 0.855 (xNa ₂ O-yNb ₂ O y = 6 y = 12.5 P/Nb 6.715 3 Eu 3.90 3.76 La 3.85 3.78	PTION THRESHOLD OF $O^{z^{-}} \rightarrow Nb^{3^{+}}$ ELECTRONIC TR. COMPOUNDS AND IN THE INVESTIGA Glass 0.855 ($xNa_2O-yNb_2O_5-zP_2O_5$)-0.14 y = 6 $y = 12.5$ $y = 18P/Nb 6.715 3 1.953Eu 3.90 3.76 3.65La 3.85 3.78 3.57$	PTION THRESHOLD OF $O^{2-} \rightarrow Nb^{3+}$ ELECTRONIC TRANSFER IN NI COMPOUNDS AND IN THE INVESTIGATED GLASSES Glass 0.855 (xNa ₂ O-yNb ₂ O ₅ -zP ₂ O ₅)-0.145 Ln ₂ O ₃ y = 6 $y = 12.5$ $y = 18$ $y = 24P/Nb 6.715 3 1.953 1.358Eu 3.90 3.76 3.65 3.54La 3.85 3.78 3.57 3.42$	PTION THRESHOLD OF $O^{2-} \rightarrow N0^{3+}$ ELECTRONIC TRANSFER IN NIOBIUM CRYST COMPOUNDS AND IN THE INVESTIGATED GLASSES Glass 0.855 (xNa ₂ O-yNb ₂ O ₅ -zP ₂ O ₅)-0.145 Ln ₂ O ₃ y = 6 $y = 12.5$ $y = 18$ $y = 24$ NbPO ₅ P/Nb 6.715 3 1.953 1.358 1 Eu 3.90 3.76 3.65 3.54 3.49 La 3.85 3.78 3.57 3.42	

TABLE I ЪT.

 α -NbPO₅ and NaNbO₃. The strong absorption observed at high energy results from the formal electronic transfer:

$$O^{2-}(2p^6) + Nb^{5+}(4d^\circ) \rightarrow O^{-}(2p^5) + Nb^{4+}(4d^1t_{2g}).$$

The values of the corresponding absorption thresholds E are given in Table I.

It is worthwhile to analyze the electronic transfer for the crystallized phases as a function of their crystal structure. The structure of NaNbO₃ (14) involves NbO₆ groups sharing their corners in the three space directions. The α -NbPO₅ lattice is built up by chains of NbO6 octahedra sharing four out of six oxygen atoms with PO₄ groups (15). The increase of the charge transfer energy E from NaNbO₃ to α -NbPO₅ results from the strengthening of the ionic character of the Nb-O bonds due to the presence of the strong covalent PO_4 groups.

As previously pointed out (under Section III) niobium is present in the glasses as NbO₆ octahedra. These groups share some of their oxygen atoms with PO₄ groups. For glasses with low niobium content the Nb-O bonds have a relatively strong ionic character due to competition with many PO4 tetrahedra, leading to a high E value. As a consequence, as niobium content increases and the relative number of PO₄ groups is lowered, E is expected to decrease. For the niobium-richest glasses the absorption threshold is close to that of α -NbPO₅, which is consistent with the presence of corner-sharing NbO₆ groups.

V. Use of Nd³⁺ and Eu³⁺ as Local Structural Probes

1. Investigation Method

The Nd³⁺ absorption of the neodymiumdoped glasses has been investigated in the scope of the Judd-Ofelt theory (16, 17).

The Judd-Ofelt parameters Ω_2 , Ω_4 , and Ω_6 are related to:

(i) the odd crystal field coefficients $A_{t,p}$ which depend on the local site symmetry, and

(ii) the $\theta(t, \lambda)$ parameter which can be considered as measuring the degree of covalency of the rare earth-ligand bond.

From the neodymium absorption spectra one can calculate the Ω_{λ} parameters. Ω_2 is the most sensitive to the neodymium-ligand bond covalency. For instance, Ω_2 decreases strongly for glasses by introduction of fluorine but increases in presence of sulfur (18).

Once the Ω_{λ} parameters have been calculated, the radiative decay probability of the ${}^{4}F_{3/2}$ level may be determined and compared with the experimental one. In order to separate the respective contributions of $A_{t,p}$ and $\theta(t, \lambda)$, the europium fluorescense has been



FIG. 4. Nd³⁺ absorption spectrum of the y = 12.5 neodymium-doped glass (thickness = 0.364 cm).

analyzed as a complementary measurement.

When Eu³⁺ occupies low symmetry sites, which occurs generally in glasses, its emission spectrum may contain electric-dipole lines such as those corresponding to ${}^{5}D_{0} \rightarrow$ ${}^{7}F_{0,2,4,6}$. But, on the other hand, the ${}^{5}D_{0} \rightarrow$ ${}^{7}F_{1}$ transition which is of magnetic dipole nature is always observed and its intensity is independent of the site symmetry. Therefore, it is generally accepted that the R = ${}^{15}D_{0} \rightarrow {}^{7}F_{2}/{}^{15}D_{0} \rightarrow {}^{7}F_{1}$ ratio represents a real assymmetry parameter for the Eu³⁺ site (19).

2. Nd³⁺ Absorption and Emission Spectra

As an example, Fig. 4 shows the absorption spectrum of the y = 12.5 glass containing $0.5 \times 10^{20} \text{ Nd}^{3+} \text{ cm}^{-3}$. The observed bands correspond to various transitions between the ground state ${}^{4}I_{9/2}$ and excited levels. The Ω_{λ} parameters calculated for the various glasses are listed in Table II. The variation of Ω_2 as a function of the Nb₂O₅ content is given in Fig. 5. Ω_2 increases with higher niobium concentration reaching 6.8 $\pm 0.3 \times 10^{-20}$ cm² for the glass corresponding to y = 24. Note that this value is one of the highest so far obtained for phosphate glasses. It indicates a high covalency of the Nd-O bonds. Figure 6 shows the neodymium ${}^{4}F_{3/2} \rightarrow {}^{4}I_{11/2}$ emission spectrum under broadband excitation. The wavelength of the emission maximum and the effective bandwidth are given in Table II. The intensity of the ${}^{4}F_{3/2} \rightarrow {}^{4}I_{11/2}$ emission for both niobium-richest glasses is about twice that of the two others.

The decay curves of the ${}^{4}F_{3/2} \rightarrow {}^{4}I_{11/2}$ Nd³⁺ emission have been recorded between 50 and 1500 μ sec. They are of exponential nature whatever the glass composition. The radiative lifetime and the corresponding measured radiative desexcitation probability A_{rad} are given in Table II. These data agree fairly well with the theoretical values calculated from the Judd-Ofelt theory.

TABLE II

Radiative Properties of the Neodymium-Doped Glasses with Composition 0.855 ($xNa_2O-yNb_2O_5-zP_2O_5$)-0.145 Nd₂O₃ (x + y + z) = 100; N_{Nd³⁺} $\approx 0.5 \times 10^{20}$ Ions $\times \text{ cm}^{-3}$

	y = 6	y = 12.5	y = 18	y = 24
Judd-Ofelt intensity parameters				
$\Omega_2 \pm 0.3 \times 10^{-20} \text{ cm}^2$	4.9	5.7	6.5	6.8
$\Omega_4 \pm 0.3 \times 10^{-20} \text{ cm}^2$	5.0	4.7	4.3	4.4
$\Omega_6 \pm 0.3 \times 10^{-20} \text{ cm}^2$	5.4	4.9	4.9	4.7
⁴ F _{3/2} lifetime				
Measured: $\tau_M \pm 10 \ \mu sec$	340	321	314	292
Calc. $\tau_{\rm rad} \pm 10 \ \mu \rm sec$	352	320	302	276
Emission probability				
Measured $1/\tau_{\rm M} \pm 100 \ {\rm sec}^{-1}$	2940	3115	3185	3425
Calc. $A_{\rm rad} \pm 100 {\rm sec^{-1}}$	2845	3080	3320	3620
Emission wavelength				
$\lambda_p \pm 2 \text{ nm}$	1056	1056	1058	1058
Effective line width				
$\Delta \lambda_{eff} \pm 0.5 \text{ nm}$	23.6	24	25	25.4



FIG. 5. Variations as a function of the Nb₂O₅ content of: (a) the Ω_2 Judd parameter for Nd³⁺; (b) the Nd³⁺ radiative emission probability A_{rad} ; (c) the Nd³⁺ $4F_{3/2} \rightarrow$ $4I_{1/2}$ emission intensity, (d) the Eu³⁺ intensity ratio

$$R = \frac{{}^{I_5}D_0 \rightarrow {}^7F_2}{{}^{I_5}D_0 \rightarrow {}^7F_1}$$

3. Eu³⁺ Emission Spectra

As an example, Fig. 7 gives the Eu³⁺ ${}^{5}D_{0} \rightarrow {}^{7}F_{J}$ (J = 0, 1, 2) emission spectrum under broad excitation recorded at 300 K for the y = 12.5 glass. The intensity of the ${}^{5}D_{0} \rightarrow {}^{7}F_{2}$ transition is stronger than that of the two other lines. Figure 5 gives the variation of the intensity ratio R of Eu³⁺ with Nb₂O₅ content. For low y values R is about 3.5



FIG. 6. Spectra of the Nd³⁺ ${}^{4}F_{3/2} \rightarrow {}^{4}I_{11/2}$ emission for the neodymium-doped glasses (x/z = 1.33; y = 6, 12.5, 18, and 24).



FIG. 7. ${}^{5}D_{0} \rightarrow {}^{7}F_{J} (J = 0, 1, 2)$ emission spectrum for the 0.855 (50 Na₂O-12.5Nb₂O₅-37.5P₂O₅)-0.145 Eu₂ O₃ glass (broad excitation at 392 nm; T = 300 K).

whereas it reaches 5.5 for the niobium-rich glasses (y = 18 and y = 24).

4. Discussion

The substitution of PO₄ tetrahedra for NbO₆ octahedra tends to strengthen the rare earth-oxygen bond in the glass lattice. This evolution is clearly illustrated by the increase of Ω_2 with niobium content. The variation of the Nd³⁺ ${}^{4}F_{3/2} \rightarrow {}^{4}I_{11/2}$ emission intensity illustrates the existence of two types of glasses: those with a low niobium oxide content (y < 12.5) and those with a higher one (y = 18 and 24).

The *R* variation deduced from the europium emission spectra supports the same conclusion: a significant increase of the assymmetry of the europium site is observed when y rises from 12.5 to 18. This evolution can be explained by the replacement of almost regular PO_4 groups by more or less distorted NbO₆ octahedra which may modify the oxygen surrounding of the rare earth.

VI. Conclusions

The structural evolution of glasses with molar compositions $0.855 (xNa_2O-yNb_2O_5-zP_2O_5)-0.145 Ln_2O_3 (Ln = La, Nd, Eu, and$ <math>x + y + z = 100) can be understood by considering the material having a medium composition (y = 12.5). It has been shown that the network former of such a glass seems to keep the Nasicon structure in which all oxygen atoms are common between NbO_6 and PO_4 polyhedra. In the niobium-poor glass (y = 6) corner-sharing NbO₆ groups are inserted within modified metaphosphate chains. As recrystallization of this glass gives rise to formation of Na₄ $Nb(PO_4)_3$, a partial demixion cannot be excluded. On the other hand, the network former of the niobium-rich glasses appears to be made up of chains containing NbO₆ octahedra sharing corners. As a consequence the cages of the covalent skeleton accommodating rare-earth ions appear to be more and more distorted at higher niobium concentrations.

References

- 1. J. H. DAYANI, Dissertation, Vanderbilt University, Nashville, TN (1977).
- A. M. GLASS, K. NASSAU, AND J. T. NEGRAN, J. Appl. Phys. 49, 1075 (1978).
- A. M. GLASS, M. E. LINES, K. NASSAU, AND J. W. SHIEVER, Appl. Phys. Lett. 31(4), 249 (1977).
- N. G. ATANOV, I. M. BOUGNISKI, E. N. KO-RIAGINA, Y. N. KRACIPOV, Y. A. POLIAKOV, A. F. SOLOXA, Y. V. TSAPKI, AND G. V. ELLERT,

Izv. Akad. Nauk SSSR, Neorg. Mater. 10, 909 (1974).

- 5. A. EL JAZOULI, R. BROCHU, J. C. VIALA, R. OLAZCUAGA, C. DELMAS, AND G. LE FLEM, Ann. Chim. (Paris) 7, 285 (1982).
- A. EL JAZOULI, C. PARENT, J. M. DANCE, AND G. LE FLEM, C.R. Acad. Sci. Paris 303, 1005 (1986).
- V. FAWCETT, D. A. LONG, AND L. H. TAYLOR, *in* "Proceedings of the 5th International Conference on Raman Spectroscopy" (H. F. Schulz, Ed.), Freiburg (1976).
- VU TI BICH, M. PROD'HOMME, M. JOUAN, AND NGUYEN QUY DAO, Verres Refract. 39(1), 13 (1985).
- 9. E. HUSSON AND Y. REPELIN, C.R. Acad. Sci. Ser. 2 296, 105 (1983).
- G. T. STRANFORD AND R. A. CONDRATE, J. Solid State Chem. 52, 248 (1984).
- 11. G. T. STRANFORD AND R. A. CONDRATE, J. Mater. Sci. Lett. 3, 303 (1984).
- G. BURNS, T. D. AXE, AND D. F. O'KANE, Solid State Commun. 7, 933 (1969).
- 13. W. ROTH, J. Res. Natl. Bur. Stand. 62, 27 (1959).
- 14. H. D. MEGAW, Acta Crystallogr. Sect. A 24, 589 (1968).
- 15. J. M. LONGO AND P. K. KIERKEGAARD, Acta Chem. Scand. 20, 72 (1966).
- 16. B. R. JUDD, Phys. Rev. 127, 750 (1962).
- 17. G. S. OFELT, J. Chem. Phys. 37, 511 (1962).
- M. J. WEBER, in "Proceedings of the International Conference on Lasers," New Orleans (1982).
- 19. C. LINARES, M. BLANCHARD, AND F. GAUME-MAHN, in "Proceedings of the 7th Rare Earth Symposium," Nauka, Moscow (1972).